

Flux Integrals: Stokes' and Gauss' Theorems

Ryan C. Daileda

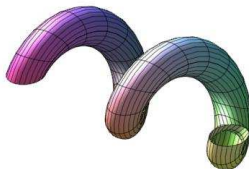
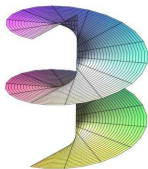
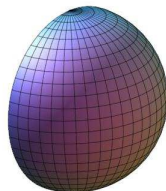
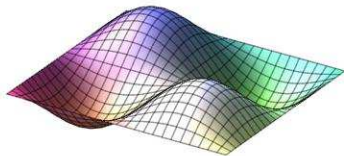


Trinity University

Calculus III
December 4, 2012

Surfaces

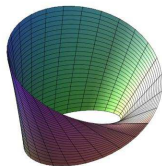
A *surface* S is a subset of \mathbb{R}^3 that is “locally planar,” i.e. when we zoom in on any point $P \in S$, S looks like a piece of a plane.



Orientable surfaces

A surface S is *orientable* if it is “two sided.”

- Every surface shown above is orientable.
- The *Möbius band* is *not* orientable.



If S is an oriented surface, an *orientation* of S is a choice of a particular side of S as “positive.”

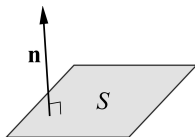
Planar flux

If S is an oriented (finite) part of a plane and $\mathbf{F} = a\mathbf{i} + b\mathbf{j} + c\mathbf{k}$ is a constant vector field, the *flux of \mathbf{F} through S* is defined to be

$$\text{comp}_{\mathbf{n}}(\mathbf{F}) A(S)$$

where:

- \mathbf{n} is the normal vector to S , in the “positive” direction;
- $A(S)$ is the area of S .



If F represents the “flow” of some quantity, then the flux is the amount of “stuff” that passes through S in one unit of time.

General flux

Suppose S is a more general oriented surface, and $\mathbf{F} = \mathbf{F}(x, y, z)$ is a possibly nonconstant vector field.

- Subdivide S into (approximately) planar pieces with (inherited) normal vectors \mathbf{n}_j and areas ΔS_j .
- Choose a point P_j in the j th subdivision, and assume that $\mathbf{F} \approx \mathbf{F}(P_j)$ throughout this subdivision.
- Compute the “local planar flux” on each subdivision and add to get the total approximate flux:

$$\sum_j \text{comp}_{\mathbf{n}_j}(\mathbf{F}(P_j)) \Delta S_j.$$

- Take the limit as the areas of the subdivisions tend to zero to get the *flux of \mathbf{F} through S* :

$$\iint_S \mathbf{F} \cdot d\mathbf{S} = \lim_{\Delta S \rightarrow 0} \sum_j \text{comp}_{\mathbf{n}_j}(\mathbf{F}(P_j)) \Delta S_j.$$

Remarks

- If we assume that \mathbf{n}_j is a unit vector, then

$$\begin{aligned}\text{comp}_{\mathbf{n}_j}(\mathbf{F}(P_j)) \Delta S_j &= (\mathbf{F}(P_j) \cdot \mathbf{n}_j) \Delta S_j \\ &= \mathbf{F}(P_j) \cdot (\Delta S_j \mathbf{n}_j) \\ &= \mathbf{F}(P_j) \cdot \Delta \mathbf{S}_j,\end{aligned}$$

where $\Delta \mathbf{S}_j$ is a normal vector with area ΔS_j .

The $d\mathbf{S}$ is thus meant to represent an “infinitesimal area normal vector” to S .

- As with planar flux, if \mathbf{F} represents the “flow” of some quantity, then $\iint_S \mathbf{F} \cdot d\mathbf{S}$ represents the amount of “stuff” that passes through S in one unit of time.

Computing flux integrals

In order to compute $\iint_S \mathbf{F} \cdot d\mathbf{S}$, one must first parametrize S via a two-variable vector function:

$$\mathbf{r}(u, v) = x(u, v)\mathbf{i} + y(u, v)\mathbf{j} + z(u, v)\mathbf{k}, \quad (u, v) \in D \subset \mathbb{R}^2.$$

If we define

$$\begin{aligned}\mathbf{T}_u &= \frac{\partial \mathbf{r}}{\partial u} = \frac{\partial x}{\partial u}\mathbf{i} + \frac{\partial y}{\partial u}\mathbf{j} + \frac{\partial z}{\partial u}\mathbf{k}, \\ \mathbf{T}_v &= \frac{\partial \mathbf{r}}{\partial v} = \frac{\partial x}{\partial v}\mathbf{i} + \frac{\partial y}{\partial v}\mathbf{j} + \frac{\partial z}{\partial v}\mathbf{k}\end{aligned}$$

then $\mathbf{T}_u \times \mathbf{T}_v$ is normal to S at every point. If the direction of \mathbf{n} agrees with the orientation of S , a Riemann sum argument shows

$$\iint_S \mathbf{F} \cdot d\mathbf{S} = \iint_D \mathbf{F}(\mathbf{r}(u, v)) \cdot (\mathbf{T}_u \times \mathbf{T}_v) du dv.$$

Example

Find the flux of the vector field

$$\mathbf{F} = x\mathbf{i} - z\mathbf{j} + y\mathbf{k}$$

through the portion of the sphere $x^2 + y^2 + z^2 = 4$ in the first octant, oriented toward the origin.

The portion of the sphere in question can be parametrized as

$$\begin{aligned}\mathbf{r}(u, v) &= 2 \sin u \cos v \mathbf{i} + 2 \sin u \sin v \mathbf{j} + 2 \cos u \mathbf{k}, \\ 0 &\leq u \leq \pi/2, \quad 0 \leq v \leq \pi/2.\end{aligned}$$

The tangent vectors are

$$\begin{aligned}\mathbf{T}_u &= 2 \cos u \cos v \mathbf{i} + 2 \cos u \sin v \mathbf{j} - 2 \sin u \mathbf{k}, \\ \mathbf{T}_v &= -2 \sin u \sin v \mathbf{i} + 2 \sin u \cos v \mathbf{j}.\end{aligned}$$

We have

$$\mathbf{T}_u \times \mathbf{T}_v = 4 \sin^2 u \cos v \mathbf{i} + 4 \sin^2 u \sin v \mathbf{j} + 4 \sin u \cos u \mathbf{k}$$

and

$$\mathbf{F}(\mathbf{r}(u, v)) = 2 \sin u \cos v \mathbf{i} - 2 \cos u \mathbf{j} + 2 \sin u \sin v \mathbf{k},$$

so that

$$\mathbf{F} \cdot (\mathbf{T}_u \times \mathbf{T}_v) = 8 \sin^3 u \cos^2 v.$$

Since $\mathbf{T}_u \times \mathbf{T}_v$ is oriented *outward* we have

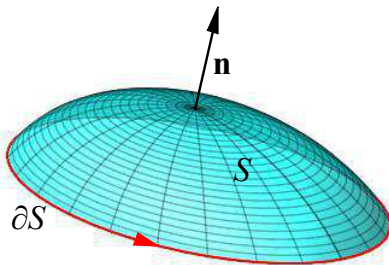
$$\iint_S \mathbf{F} \cdot d\mathbf{S} = - \int_0^{\pi/2} \int_0^{\pi/2} 8 \sin^3 u \cos^2 v \, du \, dv = -\frac{4\pi}{3}$$

A relationship between surface and line integrals

Stokes' Theorem

Let S be an oriented surface bounded by a closed curve ∂S . If \mathbf{F} is a C^1 vector field and ∂S is oriented positively relative to S , then

$$\iint_S \nabla \times \mathbf{F} \cdot d\mathbf{S} = \int_{\partial S} \mathbf{F} \cdot d\mathbf{r}.$$



Remarks

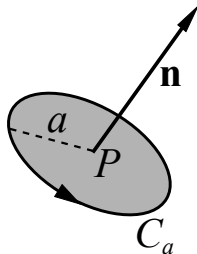
- Stokes' Theorem is another generalization of FTOC. It relates the integral of “the derivative” of \mathbf{F} on S to the integral of \mathbf{F} itself on the boundary of S .
- If $D \subset \mathbb{R}^2$ is a 2D region (oriented upward) and $\mathbf{F} = P\mathbf{i} + Q\mathbf{j}$ is a 2D vector field, one can show that

$$\iint_D \nabla \times \mathbf{F} \cdot d\mathbf{S} = \iint_D \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} dA.$$

That is, Stokes' Theorem *includes* Green's Theorem as a special case.

Interpreting the curl

Let \mathbf{F} be a vector field. Fix a point $P \in \mathbb{R}^3$ and a unit vector \mathbf{n} based at P . Let C_a denote a circle of radius a , centered at P , in the plane normal to \mathbf{n} , oriented using the right hand rule.



The tendency of \mathbf{F} to “circulate” about \mathbf{n} (in the positive sense) can be measured by

$$\lim_{a \rightarrow 0^+} \frac{1}{\pi a^2} \int_{C_a} \mathbf{F} \cdot d\mathbf{r}.$$

If D_a is the disk bounded by C_a , Stokes' Theorem implies

$$\begin{aligned}\lim_{a \rightarrow 0^+} \frac{1}{\pi a^2} \int_{C_a} \mathbf{F} \cdot d\mathbf{r} &= \lim_{a \rightarrow 0^+} \frac{1}{\pi a^2} \iint_{D_a} \nabla \times \mathbf{F} \cdot d\mathbf{S} \\ &= \lim_{a \rightarrow 0^+} \frac{1}{\pi a^2} (\nabla \times \mathbf{F})(P_0) \cdot \mathbf{n} A(D_a) \\ &= (\nabla \times \mathbf{F})(P_0) \cdot \mathbf{n}.\end{aligned}$$

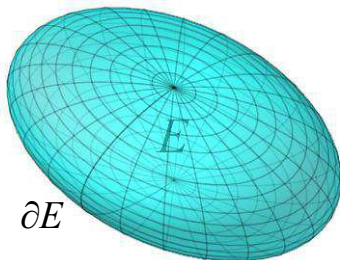
Thus, the circulation at P about \mathbf{n} is *maximized when \mathbf{n} points in the same direction as $\nabla \times \mathbf{F}$.*

A relationship between surface and triple integrals

Gauss' Theorem (a.k.a. The Divergence Theorem)

Let $E \subset \mathbb{R}^3$ be a solid region bounded by a surface ∂E . If \mathbf{F} is a C^1 vector field and ∂E is oriented outward relative to E , then

$$\iiint_E \nabla \cdot \mathbf{F} \, dV = \iint_{\partial E} \mathbf{F} \cdot d\mathbf{S}.$$



Remarks

- This can be viewed as yet another generalization of FTOC.
- Gauss' Theorem reduces computing the flux of a vector field through a *closed surface* to integrating its divergence over the region contained by that surface.
- As above, this can be used to derive a physical interpretation of $\nabla \cdot \mathbf{F}$:

$$\begin{aligned}(\nabla \cdot \mathbf{F})(P) &= \lim_{a \rightarrow 0^+} \frac{1}{V(B_a)} \iiint_{B_a} \nabla \cdot \mathbf{F} \, dV \\ &= \lim_{a \rightarrow 0^+} \frac{1}{V(B_a)} \iint_{S_a} \mathbf{F} \cdot d\mathbf{S},\end{aligned}$$

where $P \in \mathbb{R}^3$, B_a and S_a are the solid ball and sphere (respectively) of radius a centered at P .

A vast generalization

- We have studied various types of differentiation and integration in 2 and 3 dimensions.
- These can be generalized to arbitrary dimension n using the notions of “manifold” and “differential form.”
- The following theorem unifies and extends much of our integration theory in one statement.

Generalized Stokes Theorem

If M is an n -dimensional “manifold with boundary,” and ω is a “differential $(n - 1)$ -form,” then

$$\int_M d\omega = \int_{\partial M} \omega.$$