

# Vector Fields

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Calculus III

# Introduction

Today we will introduce another kind of vector-valued function, namely vector fields.

Vector fields and their line integrals are important in applications.

A vector field is simply a “repackaged” differential 1-form, and this is where we will begin.

# Differential 1-Forms

## Definition

An expression of the form

$$\omega = P(x, y) dx + Q(x, y) dy = P dx + Q dy$$

is called a (*differential*) 1-form.

The *domain* of a 1-form  $\omega$  is the intersection of the domains of  $P$  and  $Q$ .

We say that  $\omega$  is continuous, differentiable, etc. provided  $P$  and  $Q$  are both continuous, differentiable, etc.

From now on we assume  $\omega$  is always  $C^1$ , i.e. the first order partial derivatives of  $P$  and  $Q$  exist and are continuous.

# Integration of 1-Forms

Differential 1-forms associate numbers to oriented curves.

Specifically, every 1-form  $\omega$  with domain  $D$  yields a linear function

$$E_\omega : \{\text{oriented curves in } D\} \rightarrow \mathbb{R},$$

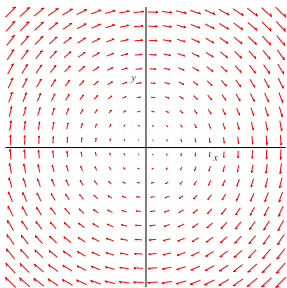
defined by  $E_\omega(C) = \int_C \omega$ .

Suppose  $\omega = P dx + Q dy$  and  $C$  is given by  $\mathbf{r}(t) = \langle x(t), y(t) \rangle$  with  $a \leq t \leq b$ . Then by definition

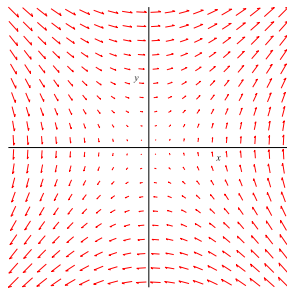
$$\begin{aligned} \int_C \omega &= \int_a^b P(x(t), y(t)) x'(t) + Q(x(t), y(t)) y'(t) dt \\ &= \int_a^b \underbrace{\langle P, Q \rangle}_{\mathbf{F}} \cdot \underbrace{\mathbf{r}'(t) dt}_{d\mathbf{r}} = \int_C \mathbf{F} \cdot d\mathbf{r}. \end{aligned}$$

The vector-valued function  $\mathbf{F}(x, y) = \langle P(x, y), Q(x, y) \rangle = P\mathbf{i} + Q\mathbf{j}$  is called a *vector field*.

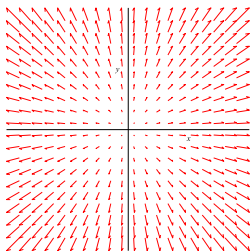
We visualize a vector field  $\mathbf{F}$  by drawing the vector  $\mathbf{F}(x, y)$  with its tail at  $(x, y)$ .



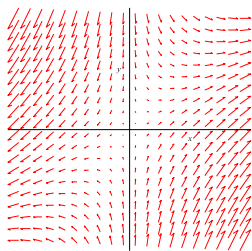
$$\mathbf{F} = \langle y, -x \rangle$$



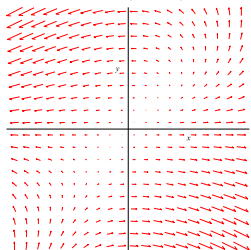
$$\mathbf{F} = \langle y, x \rangle$$



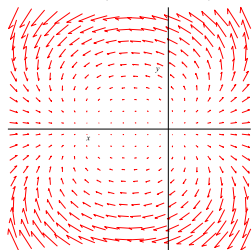
$$\mathbf{F} = \langle x, y \rangle$$



$$\mathbf{F} = \langle x, x - y \rangle$$



$$\mathbf{F} = \langle x - y, xy \rangle$$



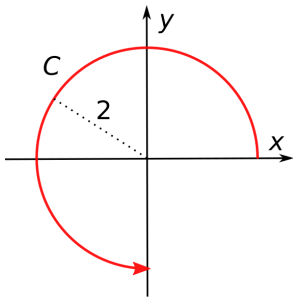
$$\mathbf{F} = \langle x^2 + x - y^2, 2xy + y \rangle$$

# Example

## Example 1

If  $\mathbf{F} = \langle x - y, xy \rangle$ , compute  $\int_C \mathbf{F} \cdot d\mathbf{r}$ , where  $C$  is  $x^2 + y^2 = 4$  from  $(2, 0)$  to  $(0, -2)$ , counterclockwise.

*Solution.*



We parametrize  $C$  using polar coordinates:

$$\mathbf{r}(t) = \langle 2 \cos t, 2 \sin t \rangle,$$

with  $0 \leq t \leq 3\pi/2$ .

We have  $x = 2 \cos t, y = 2 \sin t$  and  $\mathbf{r}'(t) = \langle -2 \sin t, 2 \cos t \rangle$ , so that

$$\begin{aligned} & \int_C \mathbf{F} \cdot d\mathbf{r} \\ &= \int_0^{3\pi/2} \langle 2 \cos t - 2 \sin t, (2 \cos t)(2 \sin t) \rangle \cdot \langle -2 \sin t, 2 \cos t \rangle dt \\ &= \int_0^{3\pi/2} -4 \sin t \cos t + 4 \sin^2 t + 8 \cos^2 t \sin t dt \\ &= \int_0^{3\pi/2} -4 \sin t \cos t + 2(1 - \cos 2t) + 8 \cos^2 t \sin t dt \\ &= -2 \sin^2 t + 2t - \sin 2t - \frac{8}{3} \cos^3 t \Big|_0^{3\pi/2} \\ &= -2 + 3\pi + \frac{8}{3} = \boxed{3\pi + \frac{2}{3}}. \end{aligned}$$



## Interpreting Line Integrals

Given a vector field  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j}$  and a curve  $C$  parametrized by  $\mathbf{r}(t)$  with  $a \leq t \leq b$ , we see that

$$\begin{aligned}\int_C \mathbf{F} \cdot d\mathbf{r} &= \int_a^b \mathbf{F}(\mathbf{r}(t)) \cdot \mathbf{r}'(t) dt \\ &= \int_a^b \frac{\mathbf{F} \cdot \mathbf{r}'}{|\mathbf{r}'|} |\mathbf{r}'| dt \\ &= \int_C \text{proj}_{\mathbf{T}}(\mathbf{F}) ds,\end{aligned}$$

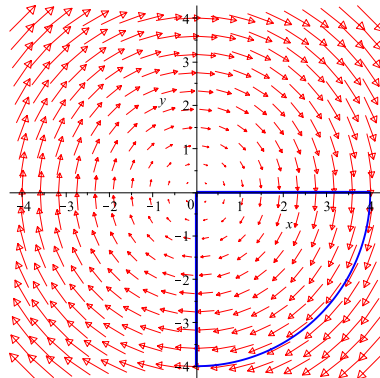
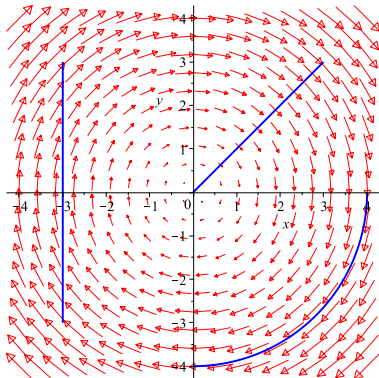
where  $\mathbf{T}(t) = \mathbf{r}'(t)/|\mathbf{r}'(t)|$  is the unit tangent vector.

So  $\int_C \mathbf{F} \cdot d\mathbf{r}$  is a net measurement of the tendency of  $\mathbf{F}(x, y)$  to “agree with” motion along  $C$ .

# Example

## Example 2

For each curve  $C$  shown, choose an orientation and determine if  $\int_C \mathbf{F} \cdot d\mathbf{r}$  is positive, negative or zero.



## A Physical Interpretation

If  $\mathbf{F}(x, y)$  represents a force at each point (e.g. gravity), and  $\mathbf{r}(t)$  is the position of a particle at time  $t$ , subject to  $\mathbf{F}$ , then

$$\mathbf{F}(\mathbf{r}(t_i)) \cdot \mathbf{r}'(t_i)\Delta t_i \approx \mathbf{F}(\mathbf{r}(t_i)) \cdot \Delta \mathbf{r}_i,$$

which is the work done by  $\mathbf{F}(\mathbf{r}(t_i))$  done through the displacement  $\Delta \mathbf{r}_i$ .

So at the level of Riemann sums we have

$$\sum_{i=1}^n \mathbf{F}(\mathbf{r}(t_i)) \cdot \mathbf{r}'(t_i)\Delta t_i \approx \sum_{i=1}^n \mathbf{F}(\mathbf{r}(t_i)) \cdot \Delta \mathbf{r}_i,$$

the RHS approximating the total work done by  $F$  as we move along  $\mathbf{r}(t)$ .

This approximation improves indefinitely as  $\Delta t \rightarrow 0$ , so that

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \text{Total work done by } \mathbf{F} \text{ along } C.$$

This suggests that we may interpret the general line integral  $\int_C \mathbf{F} \cdot d\mathbf{r}$  as a measurement of the "energy" of  $C$ , relative to  $\mathbf{F}$ .

Put another way, if  $\omega = P dx + Q dy$  is the 1-form corresponding to  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j}$ , then the "evaluation"  $E_\omega(C)$  represents the energy of  $C$  relative to  $\omega$ .

A bit more specifically, suppose  $\mathbf{F}(x, y)$  is a force field and  $\mathbf{r}(t)$  represents the motion of a particle subject *only* to the force  $\mathbf{F}$ .

Newton's Second Law gives  $\mathbf{F} = m\mathbf{r}''$  so that

$$\begin{aligned}\int_C \mathbf{F} \cdot d\mathbf{r} &= \int_a^b m\mathbf{r}''(t) \cdot \mathbf{r}'(t) dt = \frac{m}{2} \int_a^b \frac{d}{dt}(\mathbf{r}'(t) \cdot \mathbf{r}'(t)) dt \\ &= \frac{m}{2} (|\mathbf{r}'(b)|^2 - |\mathbf{r}'(a)|^2) = \frac{mv_f^2}{2} - \frac{mv_i^2}{2} = \Delta E_K\end{aligned}$$

is the change in the kinetic energy of the particle between the initial and terminal points of  $C$ .

# Conservative Vector Fields

## Definition

We call the vector field  $\mathbf{F}(x, y)$  *conservative* if  $\mathbf{F} = \nabla f$  for some function  $f(x, y)$ . In this case we call  $f$  the *potential* of  $\mathbf{F}$ .

**Warning.** In physics, the function  $-f$  is called the potential of  $\mathbf{F}$ . We will see why shortly.

**Example.** The vector field  $\mathbf{F} = y\mathbf{i} + x\mathbf{j}$  is conservative since

$$\nabla(xy) = y\mathbf{i} + x\mathbf{j} = \mathbf{F}.$$

Its potential function is  $f(x, y) = xy$ .

## Integration of Conservative Fields

Suppose that  $\mathbf{F}(x, y)$  is conservative with potential  $f(x, y)$ .

Then on any curve  $C$  we have

$$\begin{aligned}\int_C \mathbf{F} \cdot d\mathbf{r} &= \int_C f_x dx + f_y dy \\ &= \int_a^b \frac{\partial f}{\partial x} \frac{dx}{dt} + \frac{\partial f}{\partial y} \frac{dy}{dt} dt \\ &= \int_a^b \frac{d}{dt} f(\mathbf{r}(t)) dt = f(\mathbf{r}(t)) \Big|_a^b \\ &= f(\text{end of } C) - f(\text{beg. of } C).\end{aligned}$$

This is the *Fundamental Theorem of Calculus for Line Integrals*

## Theorem 1 (FTOC for Line Integrals)

If  $\mathbf{F}(x, y)$  is a conservative vector field with domain  $D$  and potential  $f(x, y)$ , then for any curve  $C$  in  $D$  from  $A$  to  $B$  we have

$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(B) - f(A).$$

Consequently:

- $\int_C \mathbf{F} \cdot d\mathbf{r}$  is path independent; it depends only on the endpoints of  $C$ .
- $\int_C \mathbf{F} \cdot d\mathbf{r} = 0$  for any closed curve  $C$  (when  $A = B$ ).



## Exact 1-Forms

Given a function  $f(x, y)$ , its *differential* is

$$df = f_x dx + f_y dy.$$

If  $\omega = P dx + Q dy$  and  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j}$  is the corresponding vector field, then

$$\mathbf{F} = \nabla f \Leftrightarrow \omega = df.$$

In this case we say that  $\omega$  is *exact*, and the FTOC for line integrals becomes

$$\int_C df = f(\text{end of } C) - f(\text{beg. of } C).$$

## Examples

### Example 3

Evaluate  $\int_C (x + 2xy) dx + x^2 dy$ , where  $C$  is the curve consisting of the line segments from  $(0, 0)$  to  $(2, 1)$  to  $(4, 3)$  to  $(5, 0)$ .

*Solution.* Notice that if  $f(x, y) = x^2y + \frac{x^2}{2}$ , then

$$df = (2xy + x) dx + x^2 dy,$$

so that FTOC implies

$$\int_C (x + 2xy) dx + x^2 dy = \int_C df = f(5, 0) - f(0, 0) = \boxed{\frac{25}{2}}.$$



### Example 4

Evaluate  $\int_C \mathbf{F} \cdot d\mathbf{r}$ , where  $\mathbf{F} = x\mathbf{i} + (y + 2)\mathbf{j}$  and  $C$  is given by  $\mathbf{r}(t) = \langle t - \sin t, 1 - \cos t \rangle$ ,  $0 \leq t \leq 2\pi$ .

*Solution.* Notice that

$$\nabla \left( \underbrace{\frac{x^2}{2} + \frac{y^2}{2} + 2y}_{f(x,y)} \right) = x\mathbf{i} + (y + 2)\mathbf{j} = \mathbf{F}.$$

Therefore, by FTOC,

$$\int_C \mathbf{F} \cdot d\mathbf{r} = f(\mathbf{r}(2\pi)) - f(\mathbf{r}(0)) = f(2\pi, 0) - f(0, 0) = \boxed{2\pi^2}.$$

## Remarks

Let  $\mathbf{F} = P\mathbf{i} + Q\mathbf{j}$  and  $\omega = P dx + Q dy$ .

Given  $f(x, y)$ , we can view  $\nabla f$  and  $df$  as different versions of the *(total) derivative* of  $f$ .

Conversely, if  $\mathbf{F}/\omega$  is conservative/exact, we can view the potential  $f(x, y)$  as the antiderivative.

FTOC then says: if  $\mathbf{F}/\omega$  has an antiderivative  $f(x, y)$ , then

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_C \omega = f(x, y) \Big|_{\text{beg.}}^{\text{end}},$$

which is an exact analogue of the single variable FTOC.

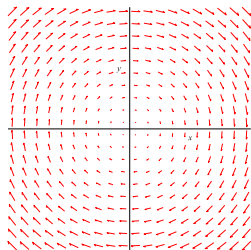
## Warning!

Unlike the single variable FTC, however, *not every vector field/1-form has an antiderivative!*

### Example 5

Use the graph of  $\mathbf{F}(x, y) = y\mathbf{i} - x\mathbf{j}$  to show that it is not conservative.

*Solution.* Here's the graph of  $\mathbf{F}$ :



If  $C$  is any clockwise circle centered at the origin, then

$$\int_C \mathbf{F} \cdot d\mathbf{r} > 0.$$



## Conservation of Energy

Suppose that  $\mathbf{F} = \nabla f$  is a conservative force field.

If a particle travels from  $A$  to  $B$  along a curve  $C$  subject only to  $\mathbf{F}$ , we have

$$E_K(B) - E_K(A) = \Delta E_K = \int_C \mathbf{F} \cdot d\mathbf{r} = f(B) - f(A),$$

or

$$E_K(B) - f(B) = E_K(A) - f(A).$$

If we define the *potential energy* to be  $E_P = -f$ , this becomes

$$E_K(B) + E_P(B) = E_K(A) + E_P(A),$$

which is the *law of conservation of energy*.

## Forms and Fields in $\mathbb{R}^3$

A 1-form in 3 variables looks like

$$\omega = P(x, y, z) dx + Q(x, y, z) dy + R(x, y, z) dz = P dx + Q dy + R dz,$$

which corresponds to the 3D vector field

$$\mathbf{F}(x, y, z) = \langle P(x, y, z), Q(x, y, z), R(x, y, z) \rangle = P\mathbf{i} + Q\mathbf{j} + R\mathbf{k}.$$

We say that  $\mathbf{F}/\omega$  is conservative/exact provided

$$\mathbf{F} = \nabla f \quad \text{or} \quad \omega = df = f_x dx + f_y dy + f_z dz.$$

The 3D analogue of the FTC for line integrals also holds.

## Example

### Example 6

Evaluate  $\int_C \mathbf{F} \cdot d\mathbf{r}$ , where  $\mathbf{F} = (y + z)\mathbf{i} + (x + z)\mathbf{j} + (x + y)\mathbf{k}$  and  $C$  is the line segment from  $(1, 0, 0)$  to  $(3, 4, 2)$ .

*Solution.* Observe that

$$\nabla \underbrace{(xy + xz + yz)}_f = (y + z)\mathbf{i} + (x + z)\mathbf{j} + (x + y)\mathbf{k}.$$

Therefore

$$\int_C \mathbf{F} \cdot d\mathbf{r} = \int_C \nabla f \cdot d\mathbf{r} = f(3, 4, 2) - f(1, 0, 0) = \boxed{26}.$$