

More Examples of the Circular Membrane

R. C. Daileda



Trinity University

Partial Differential Equations

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Recall

The shape of a vibrating circular membrane of radius a is given by

$$\begin{aligned}
 u(r, \theta, t) = & \sum_{n=1}^{\infty} a_{0n} J_0(\lambda_{0n} r) \cos(c \lambda_{0n} t) \\
 & + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} J_m(\lambda_{mn} r) (a_{mn} \cos(m\theta) + b_{mn} \sin(m\theta)) \cos(c \lambda_{mn} t) \\
 & + \sum_{n=1}^{\infty} a_{0n}^* J_0(\lambda_{0n} r) \sin(c \lambda_{0n} t) \\
 & + \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} J_m(\lambda_{mn} r) (a_{mn}^* \cos(m\theta) + b_{mn}^* \sin(m\theta)) \sin(c \lambda_{mn} t)
 \end{aligned}$$

where...

Recall (continued)

- $\lambda_{mn} = \frac{\alpha_{mn}}{a}$ and α_{mn} is the n th positive zero of $J_m(x)$.
- a_{mn} , b_{mn} (resp. a_{mn}^* , b_{mn}^*) are given by integrals involving the initial shape (resp. initial velocity), e.g. when $m \neq 0$

$$a_{mn} = \frac{\overbrace{1}^{\text{1 when } m=0} \cdot 2}{\underbrace{\pi a^2}_{\pi c \alpha_{mn} a \text{ for } *} J_{m+1}^2(\alpha_{mn})} \int_0^{2\pi} \int_0^a \overbrace{f(r, \theta)}^{g(r, \theta) \text{ for } *} \underbrace{\cos(m\theta)}_{\text{sine for } b_{mn}} J_m(\lambda_{mn} r) r \, dr \, d\theta.$$

- $a_{mn} = b_{mn} = 0$ (resp. $a_{mn}^* = b_{mn}^* = 0$) for $m \neq 0$ when f (resp. g) is *radially symmetric*, i.e. $f(r, \theta) = f(r)$.

A non-symmetric example

Example

Solve the vibrating membrane problem with $a = c = 1$ and initial conditions

$$f(r, \theta) = r(1 - r^4) \cos \theta, \quad g(r, \theta) = 0.$$

Since $g \equiv 0$, $a_{mn}^* = b_{mn}^* = 0$ for all m, n . We also have

$$\begin{aligned} b_{mn} &= \frac{2}{\pi J_{m+1}^2(\alpha_{mn})} \int_0^{2\pi} \int_0^1 r(1 - r^4) \cos \theta J_m(\alpha_{mn}r) \sin(m\theta) r \, dr d\theta \\ &= \frac{2}{\pi J_{m+1}^2(\alpha_{mn})} \underbrace{\int_0^{2\pi} \cos \theta \sin(m\theta) \, d\theta}_0 \int_0^1 r(1 - r^4) J_m(\alpha_{mn}r) r \, dr \\ &= 0 \quad \text{for all } m, n. \end{aligned}$$

Additionally,

$$\begin{aligned}
 a_{0n} &= \frac{1}{\pi J_1^2(\alpha_{0n})} \int_0^{2\pi} \int_0^1 r(1-r^4) \cos \theta J_0(\alpha_{0n}r) r \, dr d\theta \\
 &= \frac{1}{\pi J_1^2(\alpha_{0n})} \underbrace{\int_0^{2\pi} \cos \theta \, d\theta}_0 \int_0^1 r(1-r^4) J_0(\alpha_{0n}r) r \, dr \\
 &= 0,
 \end{aligned}$$

and

$$\begin{aligned}
 a_{mn} &= \frac{2}{\pi J_{m+1}^2(\alpha_{mn})} \int_0^{2\pi} \int_0^1 r(1-r^4) \cos \theta J_m(\alpha_{mn}r) \cos(m\theta) r \, dr d\theta \\
 &= \frac{2}{\pi J_{m+1}^2(\alpha_{mn})} \underbrace{\int_0^{2\pi} \cos \theta \cos(m\theta) \, d\theta}_A \int_0^1 r(1-r^4) J_m(\alpha_{mn}r) r \, dr.
 \end{aligned}$$

The integral A is zero unless $m = 1$, in which case it's equal to π .
In this case

$$\begin{aligned} a_{1n} &= \frac{2}{J_2^2(\alpha_{1n})} \int_0^1 r(1-r^4)J_1(\alpha_{1n}r)r \, dr \\ &= \frac{2}{J_2^2(\alpha_{1n})} \left(\int_0^1 r^2 J_1(\alpha_{1n}r) \, dr - \int_0^1 r^6 J_1(\alpha_{1n}r) \, dr \right). \end{aligned}$$

Substituting $x = \alpha_{1n}r$ and proceeding as before one can show

$$\begin{aligned} \int_0^1 r^2 J_1(\alpha_{1n}r) \, dr &= \frac{J_2(\alpha_{1n})}{\alpha_{1n}}, \\ \int_0^1 r^6 J_1(\alpha_{1n}r) \, dr &= \frac{J_2(\alpha_{1n})}{\alpha_{1n}} - \frac{4J_3(\alpha_{1n})}{\alpha_{1n}^2} + \frac{8J_4(\alpha_{1n})}{\alpha_{1n}^3}. \end{aligned}$$

Assembling these formulae gives

$$a_{1n} = \frac{2}{J_2^2(\alpha_{1n})} \left(\frac{4J_3(\alpha_{1n})}{\alpha_{1n}^2} - \frac{8J_4(\alpha_{1n})}{\alpha_{1n}^3} \right) = \frac{8(\alpha_{1n}J_3(\alpha_{1n}) - 2J_4(\alpha_{1n}))}{\alpha_{1n}^3 J_2^2(\alpha_{1n})}.$$

Since all the other coefficients are zero,

$$u(r, \theta, t) = \cos \theta \sum_{n=1}^{\infty} \frac{8(\alpha_{1n}J_3(\alpha_{1n}) - 2J_4(\alpha_{1n}))}{\alpha_{1n}^3 J_2^2(\alpha_{1n})} J_1(\alpha_{1n}r) \cos(\alpha_{1n}t).$$

Remark: In general, one should **not** expect the solution to reduce to a single series.

A “complicated” example

Example

Solve the vibrating membrane problem with $a = 2$, $c = 1$ and initial conditions

$$f(r, \theta) = 0, \quad g(r, \theta) = r^2(2 - r) \sin^8\left(\frac{\theta}{2}\right).$$

Since $f \equiv 0$, $a_{mn} = 0$, $b_{mn} = 0$. We also have

$$b_{mn}^* = (\dots) \int_0^2 (\dots) dr \underbrace{\int_0^{2\pi} \sin^8\left(\frac{\theta}{2}\right) \sin(m\theta) d\theta}_{\text{odd, } 2\pi\text{-periodic}} = 0,$$

$$a_{0n}^* = \frac{1}{\pi \alpha_{0n} 2 J_1^2(\alpha_{0n})} \underbrace{\int_0^{2\pi} \sin^8\left(\frac{\theta}{2}\right) d\theta}_{35\pi/64 \text{ (Maple)}} \underbrace{\int_0^2 r^2(2 - r) J_0(\lambda_{0n} r) r dr}_{?},$$

and

$$a_{mn}^* = \frac{2}{\pi \alpha_{mn} 2 J_{m+1}^2(\alpha_{mn})} \underbrace{\int_0^{2\pi} \sin^8\left(\frac{\theta}{2}\right) \cos(m\theta) d\theta}_{0 \text{ if } m \geq 5 \text{ (Maple)}} \underbrace{\int_0^2 r^2(2-r) J_m(\lambda_{mn} r) r dr.}_{?}$$

The solution therefore can be written

$$u(r, \theta, t) = \sum_{m=0}^4 \sum_{n=1}^{\infty} a_{mn}^* J_m(\lambda_{mn} r) \cos(m\theta) \sin(\lambda_{mn} t),$$

although the (?) integrals are not amenable to evaluation by hand.