# Properties of unilevel block circulants

### William F. Trench<sup>\*</sup>

Trinity University, San Antonio, Texas 78212-7200, USA Mailing address: 659 Hopkinton Road, Hopkinton, NH 03229 USA

### Abstract

Let  $\mathcal{A} = \{A_0, A_1, \dots, A_{k-1}\} \subset \mathbb{C}^{d_1 \times d_2}, \zeta = e^{-2\pi i/k}, F_\ell = \sum_{m=0}^{k-1} \zeta^{\ell m} A_m, 0 \leq \ell \leq k-1, \text{ and } \mathcal{F}_A = \bigoplus_{\ell=0}^{k-1} F_\ell.$  All operations in indices are modulo k. It is well known that if  $d_1 = d_2 = 1$  then  $[A_{s-r}]_{r,s=0}^{k-1} = \Phi \mathcal{F}_A \Phi^*,$ where  $\Phi = \frac{1}{\sqrt{k}} [\zeta^{\ell m}]_{\ell,m=0}^{k-1}$ . However, to our knowledge it has not been emphasized that  $\mathcal{F}_A$  plays a fundamental role in connection with all the matrices  $[A_{s-\alpha r}]_{r,s=0}^{k-1}$ ,  $0 \le \alpha \le k-1$ , with  $d_1$ ,  $d_2$  arbitrary. We begin by adapting a theorem of Ablow and Brenner with  $d_1 = d_2 = 1$  to the case where  $d_1$  and  $d_2$  are arbitrary. We show that  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  if and only if  $A = U_{\alpha} \mathcal{F}_A P^*$  where  $U_{\alpha}$  and P are related to  $\Phi$ , P is unitary, and  $U_{\alpha}$  is invertible (in fact, unitary) if and only if  $gcd(\alpha, k) = 1$ , in which case we say that A is a proper circulant. We prove the following for proper circulants  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$ : (i)  $A^{\dagger} = [B_{r-\alpha s}]_{r,s=0}^{k-1}$  with  $B_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{\ell m} F_{\ell}^{\dagger}, \ 0 \le m \le k-1.$ (ii) Solving Az = w reduces to solving  $F_{\ell} u_{\ell} = v_{\alpha\ell}, \ 0 \le \ell \le k-1$ , where  $v_0, v_1, \ldots, v_{k-1}$  depend only on w. (iii) A singular value decomposition of A can be obtained from singular value decompositions of  $F_0, F_1, \ldots, F_{k-1}$ . (iv) The least squares problem for A reduces to independent least squares problems for  $F_0$ ,  $F_1$ , ...,  $F_{k-1}$ . (v) If  $d_1 = d_2 = d$ , the eigenvalues of  $[A_{s-r}]_{r,s=0}^{k-1}$  are the eigenvalues of  $F_0, F_1, \ldots, F_{k-1}$ , and the corresponding eigenvectors of A are easily obtainable from those of  $F_0, F_1, \ldots, F_{k-1}$ . (vi) If  $d_1 = d_2 = d$  and  $\alpha > 1$  then the eigenvalue problem for  $[A_{s-\alpha r}]_{r,s=0}^{k-1}$  reduces to eigenvalue problems for  $d \times d$  matrices related to  $F_0, F_1, \ldots, F_{k-1}$  in a manner depending upon  $\alpha$ .

### MSC: 15A09; 15A15; 15A18; 15A99

*Keywords*: Circulant; Block circulant; Cocirculant; Least squares; Discrete Fourier transform; Moore–Penrose Inverse; Eigenvalue problem; Singular value decomposition

<sup>\*</sup>e-mail:wtrench@trinity.edu

#### Introduction 1

Throughout this paper  $k \ge 2$ ,  $d_1, d_2 \ge 1$  are integers,  $\alpha \in \{0, 1, \dots, k-1\}$ , and

$$\mathbb{C}^{k:d_1 \times d_2} = \left\{ C = [C_{rs}]_{r,s=0}^{k-1} \, \big| \, C_{rs} \in \mathbb{C}^{d_1 \times d_2}, 0 \le r, s \le k-1 \right\}.$$

All arithmetic operations in indices are modulo k. We call  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \times d_2}$  an  $\alpha$ -circulant. We say that A is a proper  $\alpha$ -circulant, or simply a proper circulant, if  $gcd(\alpha, k) = 1$ . We will say that A is a standard  $\alpha$ -circulant if  $d_1 = d_2 = 1$  and denote it by  $A = [a_{s-ar}]_{r,s=0}^{k-1}$ . Of course, there is already a vast literature on standard  $\alpha$ -circulants. Matrices of the form

$$A = [A_{rs}]_{r,s=0}^{k-1} \quad \text{where} \quad A_{rs} = \begin{cases} A_{s-r}, & 0 \le r \le s \le k-1, \\ kA_{s-r}, & 0 \le s < r \le k-1, \end{cases}$$

are also called k-circulants; see e.g., [4]. We will not consider them.

We call  $[B_{r-\alpha s}]_{r,s=0}^{k-1}$  an  $\alpha$ -cocirculant, again proper if  $gcd(\alpha, k) = 1$ . This eliminates awkward terminology such as "the conjugate transpose of the Moore-Penrose inverse of an  $\alpha$ -circulant matrix is an  $\alpha$ -circulant." The Moore-Penrose inverse of an  $\alpha$ -circulant is an  $\alpha$ -cocirculant (Theorem 4).

**Remark 1** Obviously, B is an  $\alpha$ -cocirculant if and only if  $B^*$  is an  $\alpha$ -circulant. Therefore any result concerning  $\alpha$ -circulants can be applied to  $B^*$  to obtain a result concerning B.

**Remark 2** A proper  $\alpha$ -circulant  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  is also a  $\beta$ -cocirculant where  $\alpha\beta \equiv 1 \pmod{k}$ , since

$$A_{s-\alpha r} = A_{\alpha\beta s-\alpha r} = A_{-\alpha(r-\beta s)} = B_{r-\beta s}$$

with  $B_m = A_{-\alpha m}$ ,  $0 \le m \le k - 1$ . Similarly, a proper  $\beta$ -cocirculant  $B = [B_{r-\beta s}]_{r,s=0}^{k-1}$  is also an  $\alpha$ -circulant, since

$$B_{r-\beta s} = B_{\alpha\beta r-\beta s} = B_{-\beta(s-\alpha r)} = C_{s-\alpha r}$$

with  $C_m = B_{-\beta m}, 0 \le m \le k - 1.$ 

Henceforth  $\zeta = e^{-2\pi i/k}$ ,

$$E = [\delta_{\ell,m-1}]_{\ell,m=0}^{k-1}, \quad \text{and} \quad \Phi = \frac{1}{\sqrt{k}} [\zeta^{\ell m}]_{\ell,m=0}^{k-1} = \left[ \phi_0 \quad \phi_1 \quad \cdots \quad \phi_{k-1} \right]$$
(1)

(the Fourier matrix), with

$$\phi_m = \frac{1}{\sqrt{k}} \begin{bmatrix} 1\\ \zeta^m\\ \zeta^{2m}\\ \vdots\\ \zeta^{(k-1)m} \end{bmatrix}, \quad 0 \le m \le k-1.$$
(2)

It is straightforward to verify that if indices are reduced modulo k then

$$E^{p}\left([g_{\ell m}]_{\ell,m=0}^{k-1}\right)E^{-q} = [g_{\ell+p,m+q}]_{\ell,m=0}^{k-1}.$$
(3)

Setting p = 1 and q = 0 and invoking (1) yields

$$E\Phi = \frac{1}{\sqrt{k}} [\zeta^{(\ell+1)m}]_{\ell,m=0}^{k-1} = \Phi D \quad \text{with} \quad D = \text{diag}(1,\zeta,\zeta^2,\dots,\zeta^{k-1}).$$
(4)

Therefore  $E = \Phi D \Phi^*$ .

The discrete Fourier transform (DFT) of  $\{A_0, A_1, \ldots, A_{k-1}\} \subset \mathbb{C}^{d_1 \times d_2}$  is  $\{F_0, F_1, \ldots, F_{k-1}\}$  where

$$F_{\ell} = \sum_{m=0}^{k-1} \zeta^{\ell m} A_m \in \mathbb{C}^{d_1 \times d_2}, \quad 0 \le \ell \le k-1.$$
 (5)

Since  $\Phi^{-1} = \Phi^*$ ,

$$A_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{-\ell m} F_\ell, \quad 0 \le m \le k-1.$$
(6)

We denote

$$\mathcal{F}_A = \bigoplus_{\ell=0}^{k-1} F_\ell \in \mathbb{C}^{k:d_1 \times d_2}.$$
(7)

For standard circulants (5)-(7) reduce to

$$f_{\ell} = \sum_{m=0}^{k-1} a_m \zeta^{\ell m}, \quad a_m = \frac{1}{k} \sum_{\ell=0}^{k-1} f_{\ell} \zeta^{-\ell m}, \text{ and } \mathcal{F}_A = \operatorname{diag}(f_0, f_1, \dots, f_{k-1}).$$

It is well known (see, e.g., [7]) that a standard 1-circulant  $A = [a_{s-r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k \times k}$  can be written as

$$A = \Phi \mathcal{F}_A \Phi^* = \sum_{\ell=0}^{k-1} f_\ell \phi_\ell \phi_\ell^*.$$

However, to our knowledge it has not been emphasized that  $\mathcal{F}_A$  plays a fundamental role in connection with all the standard circulants  $[a_{s-\alpha r}]_{r,s=0}^{k-1}$ . (See Remark 3.)

In Section 2 we reformulate a result of Ablow and Brenner [1, Theorem 2.1] for standard  $\alpha$ -circulants to characterize  $\alpha$ -circulants in  $\mathbb{C}^{k:d_1 \times d_2}$ . We give a different characterization in Section 3:  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  if and only if  $A = U_{\alpha} \mathcal{F}_A P^*$ , where  $U_{\alpha}$  and P are related to the Fourier matrix, P is unitary, and  $U_{\alpha}$  is invertible (in fact, unitary) if and only if  $gcd(\alpha, k) = 1$ .

Since  $\mathcal{F}_A$  is independent of  $\alpha$ , some computational results concerning  $\mathcal{F}_A$  apply simultaneously to all the proper  $\alpha$ -circulants  $[A_{s-\alpha r}]_{r,s=0}^{k-1}$ . For example, in Section 4 we show that

$$A^{\dagger} = [B_{r-\alpha s}]_{r,s=0}^{k-1}$$
 where  $B_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{\ell m} F_{\ell}^{\dagger}, \quad 0 \le m \le k-1.$ 

We also prove the following for proper  $\alpha$ -circulants: (i) Solving Az = w reduces to solving  $F_{\ell}u_{\ell} = v_{\alpha\ell}, \ 0 \leq \ell \leq k-1$ , where  $v_0, v_1, \ldots, v_{k-1}$  depend only on w. (ii) A singular value decomposition of A can be obtained from singular value decompositions of  $F_0, F_1, \ldots, F_{k-1}$ . (iii) The least squares problem for A reduces to independent least squares problems for  $F_0, F_1, \ldots, F_{k-1}$ . (iv) If  $d_1 = d_2 = d$ , the eigenvalues of  $[A_{s-r}]_{r,s=0}^{k-1}$  are the eigenvalues of  $F_0, F_1, \ldots, F_{k-1}$ , and the corresponding eigenvectors of A are easily obtainable from those of  $F_0, F_1, \ldots, F_{k-1}$ . (v) If  $d_1 = d_2 = d$  and  $\alpha > 1$ , the eigenvalue problem for  $[A_{s-\alpha r}]_{r,s=0}^{k-1}$  reduces to eigenvalue problems for  $d \times d$  matrices related to  $F_0, F_1, \ldots, F_{k-1}$  in a manner depending upon  $\alpha$ .

Block circulant 1-matrices  $[A_{s-r}]_{r,s=0}^{k-1}$  have applications in preconditioning of block Toeplitz matrices; see, e.g. [8, 9].

## 2 The Ablow–Brenner theorem revisited

Recall that E and  $\Phi$  are defined in (1) and (2). Let

$$R = E \otimes I_{d_1}, \quad P_m = \phi_m \otimes I_{d_1}, \quad 0 \le m \le k - 1, \tag{8}$$

$$S = E \otimes I_{d_2}, \quad Q_m = \phi_m \otimes I_{d_2}, \quad 0 \le m \le k - 1, \tag{9}$$

$$P = \begin{bmatrix} P_0 & P_1 & \cdots & P_{k-1} \end{bmatrix}, \quad Q = \begin{bmatrix} Q_0 & Q_1 & \cdots & Q_{k-1} \end{bmatrix}, \quad (10)$$

and

$$U_{\alpha} = \left[ \begin{array}{ccc} P_0 & P_{\alpha} & P_{2\alpha} & \cdots & P_{(k-1)\alpha} \end{array} \right].$$
(11)

Since

$$P_{\ell}^* P_m = \delta_{\ell m} I_{d_1}$$
 and  $Q_{\ell}^* Q_m = \delta_{\ell m} I_{d_2}, \quad 0 \le \ell, m \le k - 1$ 

P and Q are unitary, while  $U_\alpha$  is unitary if  $\gcd(\alpha,k)=1;$  however, if  $\gcd(\alpha,k)=q>1$  and p=k/q then

$$U_{\alpha} = [\underbrace{P_0 P_{\alpha} \cdots P_{(p-1)\alpha} \cdots P_0 P_{\alpha} \cdots P_{(p-1)\alpha}}_{q}]$$

(i.e., the first p block columns are repeated q times) is not invertible. From (4) and (8)–(11),

$$RP_{\ell} = \zeta^{\ell} P_{\ell} \quad \text{and} \quad SQ_{\ell} = \zeta^{\ell} Q_{\ell}, \quad 0 \le \ell \le k - 1.$$
 (12)

Ablow and Brenner [1, Theorem 2.1] showed that  $A \in \mathbb{C}^{k \times k}$  is a standard  $\alpha$ -circulant if and only if  $EAE^{-\alpha} = A$ . We need the following adaptation of this result.

**Theorem 1** If  $A = [G_{rs}]_{r,s=0}^{k-1}$  with  $G_{rs} \in \mathbb{C}^{d_1 \times d_2}$ , then  $RAS^{-\alpha} = A$  if and only if A is an  $\alpha$ -circulant; more precisely, if and only if

$$G_{rs} = A_{s-\alpha r}, \quad 0 \le r, s \le k-1, \tag{13}$$

with

$$A_s = G_{0s}, \quad 0 \le s \le k - 1.$$
(14)

PROOF. From (3), (8), and (9),  $RAS^{-\alpha} = [G_{r+1,s+\alpha}]_{r,s=0}^{k-1}$ . Therefore we must show that (13) is equivalent to

$$G_{r+1,s+\alpha} = G_{rs}, \quad 0 \le r, s \le k-1.$$
 (15)

If (13) holds, then

$$G_{r+1,s+\alpha} = A_{(s+\alpha)-(r+1)\alpha} = A_{s-\alpha r} = G_{rs}, \quad 0 \le r, s \le k-1.$$

For the converse we must show that (14) and (15) imply (13). We prove this by finite induction on r. From (14),

$$G_{rs} = A_{s-\alpha r}, \quad 0 \le s \le k-1, \tag{16}$$

if r = 0. Suppose (16) is true for some  $r \in \{0, \ldots, k-2\}$ . Replacing s by  $s - \alpha$  in (15) and (16) yields

$$G_{r+1,s} = G_{r,s-\alpha}, \quad 0 \le r, s \le k-1,$$

and

$$G_{r,s-\alpha} = A_{s-\alpha(r+1)}, \quad 0 \le s \le k-1.$$

Therefore

$$G_{r+1,s} = A_{s-\alpha(r+1)}, \quad 0 \le s \le k-1$$

which completes the induction.  $\hfill\square$ 

Theorem 1 with  $A = B^*$  yields the following corollary.

**Corollary 1** If  $B \in \mathbb{C}^{k:d_2 \times d_1}$  then B is an  $\alpha$ -cocirculant if and only if  $S^{\alpha}BR^{-1} = B$ .

The following corollary extends [10, Corollary 1].

**Corollary 2** (i) If  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \times d_2}$  and  $B = [B_{r-\alpha s}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_2 \times d_1}$ , then  $AB = [C_{s-r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \times d_1}$  with  $C_m = \sum_{\ell=0}^{k-1} A_\ell B_{\ell-\alpha m}, 0 \le m \le k-1$ . (ii) If  $gcd(\alpha, k) = 1$  and  $\alpha\beta \equiv 1 \pmod{k}$ , then  $BA = [D_{s-r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_2 \times d_2}$ 

$$D_m = \sum_{\ell=0}^{k-1} B_\ell A_{m+\ell}, \quad 0 \le m \le k-1.$$
(17)

PROOF. (i) From Theorem 1 and Corollary 1,  $A = RAS^{-\alpha}$  and  $B = S^{\alpha}BR^{-1}$ . Therefore  $AB = RABR^{-1}$ , so Theorem 1 with R = S implies that AB is a 1-circulant. The stated formula for  $C_0, C_1, \ldots, C_{k-1}$  can be obtained by computing first block row entries of AB.

(ii) Also,  $BA = S^{\alpha}BAS^{-\alpha}$ . Applying this  $\beta$  times yields  $BA = SBAS^{-1}$ , so Theorem 1 with R = S implies that BA is a 1-circulant. Computing the first block row entries of BA yields  $D_m = \sum_{\ell=0}^{k-1} B_{-\alpha\ell}A_{m-\alpha\ell}$  and replacing  $\ell$  by  $-\beta\ell$  yields (17).  $\Box$ 

### Theorem 2 If

$$A = [A_{s-\alpha_1 r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \times d_2}$$
(18)

and

$$B = [B_{s-\alpha_2 r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_2 \otimes d_3},$$
(19)

then

$$AB = [C_{s-\alpha_1\alpha_2r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \otimes d_3},$$

$$\tag{20}$$

with

$$C_m = \sum_{\ell=0}^{k-1} A_\ell B_{m-\alpha_2\ell}, \quad 0 \le m \le k-1.$$
(21)

PROOF. Let  $R = E \otimes I_{d_1}$ ,  $S = E \otimes I_{d_2}$ , and  $T = E \otimes I_{d_3}$ . From (18), (19), and Theorem 1,

(a)  $A = RAS^{-\alpha_1}$  and (b)  $B = SBT^{-\alpha_2}$ .

Applying (b)  $\alpha_1$  times yields  $B = S^{\alpha_1}BT^{-\alpha_1\alpha_2}$ . From this and (a),  $RABT^{-\alpha_1\alpha_2} = AB$ . Now Theorem 1 implies (20), with (21) obtained by computing the entries in the first block row of AB.

Theorem 2 generalizes [1, Theorem 3.1]; namely, the product of a standard  $\alpha$ -circulant and a standard  $\beta$ -circulant is an  $\alpha\beta$ -circulant. However, [1] does not include (21).

### **3** A DFT characterization of $\alpha$ -circulants

**Theorem 3** A matrix  $A \in \mathbb{C}^{k:d_1 \times d_2}$  is an  $\alpha$ -circulant  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  if and only if it can be written as

$$A = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} Q_{\ell}^* = U_{\alpha} \mathcal{F}_A Q^*, \qquad (22)$$

where  $\{F_0, F_1, \ldots, F_{k-1}\}$  and  $\{A_0, A_1, \ldots, A_{k-1}\}$  are related as in (5) and (6) and P, Q, and  $U_{\alpha}$  are as in (8)–(11).

PROOF. Eqns. (7)–(11) imply the second equality in (22). Therefore we need only justify the first. Suppose  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  and define  $F_0, F_1, \ldots, F_{k-1}$  by (5). From (6),

$$A_{s-\alpha r} = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{-\ell(s-\alpha r)} F_{\ell}, \quad 0 \le r, s \le k-1,$$

so (8)-(11) imply that

$$A = \frac{1}{k} \sum_{\ell=0}^{k-1} \begin{bmatrix} 1 \otimes I_{d_1} \\ \zeta^{\alpha\ell} \otimes I_{d_1} \\ \vdots \\ \zeta^{(k-1)\alpha\ell} \otimes I_{d_1} \end{bmatrix} F_{\ell} \begin{bmatrix} 1 \otimes I_{d_2} \\ \zeta^{\ell} \otimes I_{d_2} \\ \vdots \\ \zeta^{(k-1)\ell} \otimes I_{d_2} \end{bmatrix}^H = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} Q_{\ell}^*.$$

For the converse, if (22) holds then (12) implies that

$$RAS^{-\alpha} = \sum_{\ell=0}^{k-1} (RP_{\alpha\ell}) F_{\ell} (S^{\alpha}Q_{\ell})^* = \sum_{\ell=0}^{k-1} (\zeta^{\alpha\ell}P_{\alpha\ell}) F_{\ell} (\zeta^{-\alpha\ell}Q_{\ell}^*) = A.$$

Therefore A is an  $\alpha$ -circulant, by Theorem 1; hence,  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  with  $A_0$ ,  $A_1, \ldots, A_{k-1}$  as in (6).  $\square$ 

**Remark 3** Theorem 3 implies that  $A \in \mathbb{C}^{k \times k}$  is a standard  $\alpha$ -circulant  $[a_{s-\alpha r}]_{r,s=0}^{k-1}$  if and only if

$$A = \Phi_{\alpha} \mathcal{F}_A \Phi^* = \sum_{\ell=0}^{k-1} f_\ell \phi_{\alpha\ell} \phi_\ell^*,$$

where  $\Phi$  is as in (1),  $\Phi_{\alpha} = \begin{bmatrix} \phi_0 & \phi_{\alpha} & \cdots & \phi_{(k-1)\alpha} \end{bmatrix}$ , and

$$f_{\ell} = \sum_{m=0}^{k-1} a_m \zeta^{\ell m}, \quad 0 \le \ell \le k-1.$$

**Corollary 3** A matrix  $B \in \mathbb{C}^{k:d_2 \times d_1}$  is an  $\alpha$ -cocirculant if and only if it can be written as  $B = \sum_{\ell=0}^{k-1} Q_\ell G_\ell P^*_{\alpha\ell}$ , where

$$G_{\ell} = \sum_{m=0}^{k-1} \zeta^{-\ell m} B_m, \quad 0 \le \ell \le k-1, \text{ and } B_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{\ell m} G_{\ell}, \quad 0 \le m \le k-1.$$

PROOF. Apply Theorem 3 to  $B^*$ .

It is well known that standard 1-circulants commute. The following corollary extends this.

**Corollary 4** Suppose  $d_1 = d_2$ ,  $gcd(\alpha, k) = 1$ , and  $\alpha\beta \equiv 1 \pmod{k}$ . Let  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$ ,  $B = [B_{s-\beta r}]_{r,s=0}^{k-1}$ ,

$$F_{\ell} = \sum_{m=0}^{k-1} \zeta^{\ell m} A_m$$
 and  $G_{\ell} = \sum_{m=0}^{k-1} \zeta^{\ell m} B_m$ .

Then AB = BA if and only if  $F_{\beta\ell}G_{\ell} = G_{\alpha\ell}F_{\ell}, 0 \leq \ell \leq k-1$ .

**PROOF.** Since  $gcd(\alpha, k) = gcd(\beta, k) = 1$ , we may change summation indices  $\ell \to \alpha \ell$  and  $\ell \to \beta \ell$ . Therefore, from Theorem 3 with Q = P,

$$A = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} P_{\ell}^* = \sum_{\ell=0}^{k-1} P_{\ell} F_{\beta\ell} P_{\beta\ell}^*, \quad B = \sum_{\ell=0}^{k-1} P_{\beta\ell} G_{\ell} P_{\ell}^* = \sum_{\ell=0}^{k-1} P_{\ell} G_{\alpha\ell} P_{\alpha\ell}^*,$$
$$AB = \sum_{\ell=0}^{k-1} P_{\ell} F_{\beta\ell} G_{\ell} P_{\ell}^*, \quad \text{and} \quad BA = \sum_{\ell=0}^{k-1} P_{\ell} G_{\alpha\ell} F_{\ell} P_{\ell}^*,$$

which implies the conclusion.  $\hfill\square$ 

# 4 Moore-Penrose inversion and singular value decomposition

Recall that the Moore-Penrose inverse  $X^\dagger$  of a matrix X is the unique matrix Y that satisfies the Penrose conditions

$$(XY)^* = XY, \quad (YX)^* = YX, \quad XYX = X, \text{ and } YXY = Y.$$

**Theorem 4** The Moore–Penrose inverse of an  $\alpha$ -circulant is an  $\alpha$ -cocirculant.

PROOF. From Theorem 1, if A is an  $\alpha$ -circulant then  $A = RAS^{-\alpha}$ . Let  $B = S^{\alpha}A^{\dagger}R^{-1}$ . We will show that A and B satisfy the Penrose conditions:

$$AB = (RAS^{-\alpha})(S^{\alpha}A^{\dagger}R^{-1}) = RAA^{\dagger}R^{*} = R(AA^{\dagger})^{*}R^{*} = (AB)^{*},$$
  
$$BA = (S^{\alpha}A^{\dagger}R^{-1})(RAS^{-\alpha}) = S^{\alpha}A^{\dagger}A(S^{\alpha})^{*} = S^{\alpha}(A^{\dagger}A)^{*}(S^{\alpha})^{*} = (BA)^{*},$$

$$ABA = (RAA^{\dagger}R^{-1})(RAS^{-\alpha}) = R(AA^{\dagger}A)S^{-\alpha} = RAS^{-\alpha} = A,$$

and

$$BAB = (S^{\alpha}A^{\dagger}AS^{-\alpha})(S^{\alpha}A^{\dagger}R^{-1}) = S^{\alpha}(A^{\dagger}AA^{\dagger})R^{-1} = S^{\alpha}A^{\dagger}R^{-1} = B.$$

Therefore  $B = A^{\dagger}$  or, equivalently,  $S^{\alpha}A^{\dagger}R^{-1} = A^{\dagger}$ . Now Corollary 1 implies that  $A^{\dagger}$  is an  $\alpha$ -cocirculant.

We can be more explicit if  $gcd(\alpha, k) = 1$ .

**Theorem 5** The Moore–Penrose inverse of a proper  $\alpha$ -circulant  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$ is the  $\alpha$ -cocirculant  $B = [B_{r-\alpha s}]_{r,s=0}^{k-1}$ , where

$$B_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{\ell m} F_{\ell}^{\dagger}, \quad 0 \le m \le k-1,$$
(23)

with

$$F_{\ell} = \sum_{m=0}^{k-1} \zeta^{\ell m} A_m, \quad 0 \le \ell \le k-1.$$

PROOF. From Theorem 3,  $A = U_{\alpha} \mathcal{F}_A Q^*$  where Q and  $U_{\alpha}$  are unitary, the latter since  $gcd(\alpha, k) = 1$ . We will first show that A and  $B = Q \mathcal{F}_A^{\dagger} U_{\alpha}^*$  satisfy the Penrose conditions:

$$AB = (U_{\alpha}\mathcal{F}_{A}Q^{*})(Q\mathcal{F}_{A}^{\dagger}U_{\alpha}^{*}) = U_{\alpha}\mathcal{F}_{A}\mathcal{F}_{A}^{\dagger}U_{\alpha}^{*} = U_{\alpha}(\mathcal{F}_{A}\mathcal{F}_{A}^{\dagger})^{*}U_{\alpha}^{*} = (AB)^{*},$$
  

$$BA = (Q\mathcal{F}_{A}^{\dagger}U_{\alpha}^{*})(U_{\alpha}\mathcal{F}_{A}Q^{*}) = Q\mathcal{F}_{A}^{\dagger}\mathcal{F}_{A}Q^{*} = Q(\mathcal{F}_{A}^{\dagger}\mathcal{F}_{A})^{*}Q^{*} = (BA)^{*},$$
  

$$ABA = (U_{\alpha}\mathcal{F}_{A}\mathcal{F}_{A}^{\dagger}U_{\alpha}^{*})(U_{\alpha}\mathcal{F}_{A}Q^{*}) = U_{\alpha}(\mathcal{F}_{A}\mathcal{F}_{A}^{\dagger}\mathcal{F}_{A})Q^{*} = U_{\alpha}\mathcal{F}_{A}Q^{*} = A,$$

and

$$BAB = (Q\mathcal{F}_A^{\dagger}\mathcal{F}_A Q^*)(Q\mathcal{F}_A^{\dagger}U_{\alpha}^*) = Q(\mathcal{F}_A^{\dagger}\mathcal{F}_A \mathcal{F}_A^{\dagger})U_{\alpha}^* = Q\mathcal{F}_A^{\dagger}U_{\alpha}^* = B.$$

Therefore

$$A^{\dagger} = B = \sum_{\ell=0}^{k-1} Q_{\ell} F_{\ell}^{\dagger} P_{\alpha\ell}^{*} = \sum_{\ell=0}^{k-1} (\phi_{\ell} \otimes I_{d_{2}}) F_{\ell}^{\dagger} (\phi_{\alpha\ell} \otimes I_{d_{1}})^{*}$$
$$= \frac{1}{k} \left[ \sum_{\ell=0}^{k-1} \zeta^{\ell(r-\alpha s)} F_{\ell}^{\dagger} \right]_{r,s=0}^{k-1} = [B_{r-\alpha s}]_{r,s=0}^{k-1},$$

from (8)-(11) and (23).

**Remark 4** Theorem 5 can also be proved by using (6) and (23) to express the entries of AB, BA, ABA, and BAB explicitly in terms of  $F_0$ ,  $F_1$ , ...,  $F_{k-1}$  and  $F_0^{\dagger}$ ,  $F_1^{\dagger}$ , ...,  $F_{k-1}^{\dagger}$ , noting that

$$\sum_{\ell=0}^{k-1} \zeta^{\ell(r-s)} = \sum_{\ell=0}^{k-1} \zeta^{\alpha\ell(r-s)} = \delta_{rs}, \quad 0 \le r, s \le k-1,$$

the latter because  $gcd(\alpha, k) = 1$ . However, this is tedious.

**Remark 5** Theorem 5 extends a result of Davis [6]: If  $A = [a_{s-r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k \times k}$ then  $A^{\dagger} = \Phi \operatorname{diag}(a_0^{\dagger}, a_1^{\dagger}, \dots, a_{k-1}^{\dagger})\Phi^*$ , where  $\Phi$  is the Fourier matrix (1),  $0^{\dagger} = 0$ , and  $a^{\dagger} = 1/a$  if  $a \neq 0$ .

**Theorem 6** Suppose  $gcd(\alpha, k) = 1$  and

$$A = [A_{s-\alpha r}]_{r,s=0}^{k-1} = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} Q_{\ell}^* = U_{\alpha} \mathcal{F}_A Q^*.$$

Let  $F_{\ell} = \Omega_{\ell} \Sigma_{\ell} \Psi_{\ell}^*$  be a singular value decomposition of  $F_{\ell}$ ,  $0 \leq \ell \leq k-1$ , and define

$$M_{\alpha} = \left[ \begin{array}{ccc} P_0 \Omega_0 & P_{\alpha} \Omega_1 & \cdots & P_{(k-1)\alpha} \Omega_{k-1} \end{array} \right]$$

and

$$N = \begin{bmatrix} Q_0 \Psi_0 & Q_1 \Psi_1 & \cdots & Q_{k-1} \Psi_{k-1} \end{bmatrix}.$$

Then

$$A = M_{\alpha} \left( \bigoplus_{\ell=0}^{k-1} \Sigma_{\ell} \right) N^*$$

is a singular value decomposition of A, except that the singular values are not necessarily ordered.

# 5 The least squares problem

Suppose  $G \in \mathbb{C}^{d_1 \times d_2}$  and consider the least squares problem for G: If  $v \in \mathbb{C}^{d_1}$ , find  $u \in \mathbb{C}^{d_2}$  such that

$$||Gu - v|| = \min_{\xi \in \mathbb{C}^{d_2}} ||G\xi - v||,$$
(24)

where  $\|\cdot\|$  is the 2-norm. This problem has a unique solution if and only if rank $(G) = d_2$ , in which case  $u = (G^*G)^{-1}G^*v$ . In any case, the optimal solution of (24) is the unique  $u_0 \in \mathbb{C}^{d_2}$  of minimum norm that satisfies (24); thus,  $u_0 = G^{\dagger}v$ . The general solution of (24) is  $u = u_0 + q$  where Gq = 0, and

$$\|Gu - v\| = \|(GG^{\dagger} - I)v\|$$

for all such u.

Now consider the following least squares problem: if  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1} \in \mathbb{C}^{k:d_1 \times d_2}$  with  $gcd(\alpha, k) = 1$  and  $w \in \mathbb{C}^{kd_1}$ , find  $z \in \mathbb{C}^{kd_2}$  such that

$$||Az - w|| = \min_{\xi \in \mathbb{C}^{kd_2}} ||A\xi - w||$$

We write

$$z = \sum_{\ell=0}^{k-1} Q_{\ell} u_{\ell} \quad \text{and} \quad w = \sum_{\ell=0}^{k-1} P_{\ell} v_{\ell} = \sum_{\ell=0}^{k-1} P_{\alpha\ell} v_{\alpha\ell}, \tag{25}$$

since substituting  $\alpha \ell$  for  $\ell$  is legitimate because  $gcd(\alpha, k) = 1$ . Since  $A = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} Q_{\ell}^*$  and  $Q_{\ell}^* Q_m = \delta_{\ell m} I_{kd_2}$ ,

$$Az - w = \sum_{\ell=0}^{k-1} P_{\alpha\ell} (F_{\ell} u_{\ell} - v_{\alpha\ell}).$$

Since  $P_{\alpha\ell}^* P_{\alpha m} = \delta_{\ell m} I_{d_1}$  (because  $gcd(\alpha, k) = 1$ ), it follows that

$$||Az - w||^2 = \sum_{\ell=0}^{k-1} ||F_{\ell}u_{\ell} - v_{\alpha\ell}||^2.$$
(26)

This implies the following theorem.

**Theorem 7** Suppose A is a proper  $\alpha$ -circulant and let z and w be as in (25). Then

$$||Az - w|| = \min_{\xi \in \mathbb{C}^{kd_2}} ||A\xi - w||$$
(27)

if and only if

$$\|F_{\ell}u_{\ell} - v_{\alpha\ell}\| = \min_{\psi_{\ell} \in C^{d_2}} \|F_{\ell}\psi_{\ell} - v_{\alpha\ell}\|, \quad 0 \le \ell \le k - 1.$$

Therefore (27) has a unique solution, given by

$$z = \sum_{\ell=0}^{k-1} Q_{\ell} (F_{\ell}^* F_{\ell})^{-1} F_{\ell}^* v_{\alpha\ell},$$

if and only rank $(F_{\ell}) = d_2, \ 0 \leq \ell \leq k-1$ . In any case, the optimal solution of (27) is

$$z_0 = \sum_{\ell=0}^{k-1} Q_\ell F_\ell^\dagger v_{\alpha\ell}$$

The general solution of (27) is  $z = z_0 + \sum_{\ell=0}^{k-1} Q_\ell u_\ell$ , where  $F_\ell u_\ell = 0, \ 0 \le \ell \le k-1$ , and

$$||Az - w||^2 = \sum_{\ell=0}^{k-1} ||(F_{\ell}F_{\ell}^{\dagger} - I_{d_1})v_{\alpha\ell}||^2$$

for all such z.

# 6 The case where $d_1 = d_2$

Throughout this section  $d_1 = d_2 = d$  and  $A = [A_{s-\alpha r}]_{r,s=0}^{k-1}$  is a proper circulant. Then (26) implies the following theorem, which reduces the problem of solving the  $kd \times kd$  system Az = w to solving k independent  $d \times d$  systems.

**Theorem 8** If A is a proper  $\alpha$ -circulant,  $z = \sum_{\ell=0}^{k-1} P_{\ell} u_{\ell}$ , and  $w = \sum_{\ell=0}^{k-1} P_{\ell} v_{\ell}$ , then Az = w if and only if

$$F_{\ell}u_{\ell} = v_{\alpha\ell}, \quad 0 \le \ell \le k - 1.$$

This and Theorem 5 imply the following theorem.

**Theorem 9** A proper  $\alpha$ -circulant

$$A = [A_{s-\alpha r}]_{r,s=0}^{k-1} = \sum_{\ell=0}^{k-1} P_{\alpha \ell} F_{\ell} P_{\ell}^*$$
(28)

is invertible if and only  $F_0, F_1, \ldots, F_{k-1}$  are all invertible. In this case

$$A^{-1} = [B_{r-\alpha s}]_{r,s=0}^{k-1} \quad with \quad B_m = \frac{1}{k} \sum_{\ell=0}^{k-1} \zeta^{\ell m} F_{\ell}^{-1}, \quad 0 \le m \le k-1,$$

and the solution of Az = w is  $z = \sum_{\ell=0}^{k-1} P_{\ell} F_{\ell}^{-1} v_{\alpha \ell}$ .

**Remark 6** Theorem 9 and Remark 2 extend [5, Theorem 1]: the inverse of a standard nonsingular  $\alpha$ -circulant is a  $\beta$ -circulant, where  $\alpha\beta \equiv 1 \pmod{k}$ . **Theorem 10** Suppose A is a proper  $\alpha$ -circulant as in (28) and  $\alpha\beta \equiv 1 \pmod{k}$ .

- (i) A is Hermitian if and only if  $P_{\beta\ell}F^*_{\beta\ell} = P_{\alpha\ell}F_{\ell}, \ 0 \le \ell \le k-1.$
- (ii) A is normal if and only if  $F_{\beta\ell}F^*_{\beta\ell} = F^*_{\ell}F_{\ell}, 0 \le \ell \le k-1$ .
- (iii) A is EP (i.e.,  $A^{\dagger}A = AA^{\dagger}$ ) if and only if  $F_{\beta\ell}F_{\beta\ell}^{\dagger} = F_{\ell}^{\dagger}F_{\ell}, 0 \leq \ell \leq k-1$ .

Proof.

From (28) and Theorem 5,

$$A = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} P_{\ell}^*, \quad A^* = \sum_{\ell=0}^{k-1} P_{\ell} F_{\ell}^* P_{\alpha\ell}^*, \quad \text{and} \quad A^{\dagger} = \sum_{\ell=0}^{k-1} P_{\ell} F_{\ell}^{\dagger} P_{\alpha\ell}^*.$$
(29)

(i) Since  $\alpha\beta \equiv 1 \pmod{k}$ , replacing  $\ell$  by  $\beta\ell$  in the second sum in (29) yields  $A^* = \sum_{\ell=0}^{k-1} P_{\beta\ell} F^*_{\beta\ell} P^*_{\ell}$ , and comparing this with the first sum in (29) yields (i). (ii) From (29),

$$AA^* = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} F_{\ell}^* P_{\alpha\ell}^* = \sum_{\ell=0}^{k-1} P_{\ell} F_{\beta\ell} F_{\beta\ell}^* P_{\ell}^* \quad \text{and} \quad A^*A = \sum_{\ell=0}^{k-1} P_{\ell} F_{\ell}^* F_{\ell} P_{\ell}^*,$$

which implies (ii).

(iii) From (29),

$$AA^{\dagger} = \sum_{\ell=0}^{k-1} P_{\alpha\ell} F_{\ell} F_{\ell}^{\dagger} P_{\alpha\ell}^{*} = \sum_{\ell=0}^{k-1} P_{\ell} F_{\beta\ell} F_{\beta\ell}^{\dagger} P_{\ell}^{*} \quad \text{and} \quad A^{\dagger}A = \sum_{\ell=0}^{k-1} P_{\ell} F_{\ell}^{\dagger} F_{\ell} P_{\ell}^{*}.$$

which implies (iii).  $\Box$ 

**Remark 7** If A is a square matrix and there is a matrix B such that ABA = A, BAB = B, and AB = BA, then B is unique and is called the group inverse of A, which is usually denoted by  $A^{\#}$ . Davis [6] noted that if  $A \in \mathbb{C}^{k \times k}$  is a standard 1-circulant then  $A^{\dagger} = A^{\#}$ . Theorem 10(iii) extends this: If  $A \in \mathbb{C}^{k:d \times d}$  is a proper  $\alpha$ -circulant and  $\alpha\beta \equiv 1 \pmod{k}$ , then  $A^{\dagger} = A^{\#}$  if and only if  $F_{\ell}^{\dagger}F_{\ell} = F_{\beta\ell}F_{\beta\ell}^{\dagger}, 0 \leq \ell \leq k-1$ .

# 7 The eigenvalue problem with $\alpha = 1$

In this section we assume that  $\alpha = 1$  and  $d_1 = d_2 = d$ . The following theorem and its proof are motivated by [2, Theorem 2].

Theorem 11 Let

$$\mathcal{S}_R = \bigcup_{\ell=0}^{k-1} \left\{ z \, \big| \, Rz = \zeta^\ell z \right\}.$$

If  $\lambda$  is an eigenvalue of A, let  $\mathcal{E}_A(\lambda)$  be the  $\lambda$ -eigenspace of A; i.e,

$$\mathcal{E}_A(\lambda) = \left\{ z \, \big| \, Az = \lambda z \right\}$$

(i) If  $\lambda$  is an eigenvalue of  $A = [A_{s-r}]_{r,s=0}^{k-1}$  then  $\mathcal{E}_A(\lambda)$  has a basis in  $\mathcal{S}_R$ .

(ii) If  $A \in \mathbb{C}^{k:d \times d}$  and has kd linearly independent eigenvectors in  $S_R$ , then A is a 1-circulant.

PROOF. (i) From Theorem 8,  $z = \sum_{\ell=0}^{k-1} P_{\ell} u_{\ell} \in \mathcal{E}_A(\lambda)$  if and only if  $F_{\ell} u_{\ell} = \lambda u_{\ell}$ ,  $0 \leq \ell \leq k-1$ . Therefore  $\lambda$  is an eigenvalue of A if and only if it is an eigenvalue of  $F_{\ell}$  for some  $\ell \in \{0, 1, \ldots, k-1\}$ . Let  $\mathcal{T}_{\lambda}$  be the subset of  $\{0, 1, \ldots, k-1\}$  for which this is true. Then  $\mathcal{E}_A(\lambda)$  consists of linear combinations of the vectors of the form  $P_{\ell}u_{\ell}$  with  $\ell \in \mathcal{T}_{\lambda}$  and  $(\lambda, u_{\ell})$  an eigenpair of  $F_{\ell}$ . Since  $RP_{\ell} = \zeta^{\ell}P_{\ell}$  (recall (12)), this completes the proof of (i).

(ii) From Theorem 1, we must show that RA = AR. If  $Az = \lambda z$  and  $Rz = \zeta^s z$  then  $RAz = \lambda Rz = \lambda \zeta^s z$  and  $ARz = \zeta^s Az = \zeta^s \lambda z$ . Hence ARz = RAz for all z in a basis for  $\mathbb{C}^{k:d \times d}$ , so AR = RA.

**Theorem 12** Let R and P be as in (8) and (10). Then the 1-circulant  $A = [A_{s-r}]_{r,s=0}^{k-1}$  is diagonalizable if and only if  $F_0, F_1, \ldots, F_{k-1}$  are all diagonalizable. In this case, if

$$F_{\ell} = T_{\ell} D_{\ell} T_{\ell}^*, \quad 0 \le \ell \le k - 1,$$

are spectral decompositions of  $F_0, F_1, \ldots, F_{k-1}$  and

$$\Psi = \left[ \begin{array}{ccc} P_0 T_0 & P_1 T_1 & \cdots & P_{k-1} T_{k-1} \end{array} \right],$$

then

$$A = \Psi\left(\bigoplus_{\ell=0}^{k-1} D_\ell\right) \Psi^*$$

is a spectral decomposition of A.

## 8 The eigenvalue problem with $\alpha > 1$

In this section we assume that  $d_1 = d_2 = d$ ,  $\alpha \in \{2, 3, \dots, k-1\}$  and,  $gcd(\alpha, k) = 1$ . From Theorem 8,  $Az = \lambda z$  if and only if  $z = \sum_{s=0}^{k-1} P_s u_s$ , where

$$F_s u_s = \lambda u_{\alpha s}, \quad 0 \le s \le k - 1. \tag{30}$$

Therefore Az = 0 if and only if  $z = \sum_{s=0}^{k-1} P_s u_s$  where  $F_s u_s = 0, 0 \le s \le k-1$ , so the makeup of the null space of A is transparent. Hence, we assume that  $\lambda \ne 0$ . Then we must consider the orbits of the permutation on  $\{0, \ldots, k-1\}$ defined by  $s \rightarrow \alpha s \pmod{k}$ . We consider an example before presenting the general discussion. Let k = 10 and  $\alpha = 3$ . The permutation of  $\{0, 1, \ldots, 9\}$  defined by  $s \to 3s \pmod{10}$  is given by

The orbits of this permutation are

$$\mathcal{O}_0 = \{0\}, \quad \mathcal{O}_1 = \{1, 3, 9, 7\}, \quad \mathcal{O}_2 = \{2, 6, 8, 4\}, \text{ and } \mathcal{O}_3 = \{5\}.$$

Therefore (30) divides into four independent systems:

(i) 
$$F_0 u_0 = \lambda u_0$$
; (ii)  $F_1 u_1 = \lambda u_3$ ,  $F_3 u_3 = \lambda u_9$ ,  $F_9 u_9 = \lambda u_7$ ,  $F_7 u_7 = \lambda u_1$ 

(iii) 
$$F_5 u_5 = \lambda u_5$$
; (iv)  $F_2 u_2 = \lambda u_6$ ,  $F_6 u_6 = \lambda u_8$ ,  $F_8 u_8 = \lambda u_4$ ,  $F_4 u_4 = \lambda u_2$ 

From (i), if  $(\lambda, u_0)$  is an eigenpair of  $F_0$  then  $(\lambda, P_0 u_0)$  is an eigenpair of A. Similarly, from(iii), if  $(\lambda, u_5)$  is an eigenpair of  $F_5$  then  $(\lambda, P_5 u_5)$  is an eigenpair of A. The analysis of (ii) and (iv) is more complicated, but identical. We will consider (ii), which is equivalent to

$$u_3 = \frac{1}{\lambda} F_1 u_1, \quad u_9 = \frac{1}{\lambda} F_3 u_3, \quad u_7 = \frac{1}{\lambda} F_9 u_9, \quad u_1 = \frac{1}{\lambda} F_7 u_7,$$
 (31)

since  $\lambda \neq 0$ . Hence,

$$u_3 = \frac{1}{\lambda} G_3 u_1 \quad u_9 = \frac{1}{\lambda^2} G_9 u_1, \quad u_7 = \frac{1}{\lambda^3} G_7 u_1, \quad \text{and} \quad u_1 = \frac{1}{\lambda^4} G_1 u_1, \quad (32)$$

where

$$G_3 = F_1, \quad G_9 = F_3 F_1, \quad G_7 = F_9 F_3 F_1, \quad \text{and} \quad G_1 = F_7 F_9 F_3 F_1.$$
 (33)

In particular, the last equalities in (32) and (33) are equivalent to  $G_1u_1 = \lambda^4 u_1$ . Therefore, if  $(\gamma, u_1)$  is an eigenpair of  $G_1$  and  $\gamma \neq 0$ , then  $\lambda = \gamma^{1/4}$  is an eigenvalue of A with the associated eigenvector

$$z = (P_1 + \gamma^{-1/4} P_3 G_3 + \gamma^{-2/4} P_9 G_9 + \gamma^{-3/4} P_7 G_7) u_1$$
  
=  $\left(P_1 + \sum_{m=1}^3 \gamma^{-m/4} P_{3^m} G_{3^m}\right) u_1.$  (34)

(Recall that subscripts are taken modulo 10.) However,  $\gamma^{1/4}e^{2\pi i r/4}$ ,  $0 \le r \le 3$ , are all fourth roots of  $\gamma$  and therefore eigenvalues of A. Replacing  $\gamma^{1/4}$  with  $\gamma^{1/4}e^{2\pi i r/4}$  in (34) shows that

$$z_r = \left(P_1 + \sum_{m=1}^3 \gamma^{-m/4} e^{-2\pi i r m/4} P_{3^m} G_{3^m}\right) u_1, \quad 0 \le r \le 3, \tag{35}$$

are the respective associated eigenvectors of A.

Now suppose the permutation  $s \to \alpha s \pmod{k}$  of  $\{0, 1, \ldots, k-1\}$  has p orbits  $\mathcal{O}_0, \mathcal{O}_1, \ldots, \mathcal{O}_{p-1}$ , and let

$$0 = s_0 < s_1 < s_2 < \dots < s_{p-1}$$
 with  $s_\ell \in \mathcal{O}_\ell$ ,  $0 \le \ell \le p-1$ .

Suppose  $\mathcal{O}_{\ell}$  has  $r_{\ell}$  distinct members; thus,

$$\mathcal{O}_{\ell} = \{s_{\ell}, \alpha s_{\ell}, \dots, \alpha^{r_{\ell}-1} s_{\ell}\} \quad \text{where} \quad \alpha^{r_{\ell}} \equiv 1 \pmod{k}, \tag{36}$$

and  $\bigcup_{\ell=0}^{p-1} \mathcal{O}_{\ell} = \{0, 1, \dots, k-1\}$ . If  $r_{\ell} = 1$  and  $(\lambda, u_{s_{\ell}})$  is an eigenpair of  $F_{s_{\ell}}$ , then  $(\lambda, P_{s_{\ell}} u_{s_{\ell}})$  is an eigenpair of A. Now consider an orbit  $\mathcal{O}_{\ell}$  with  $r_{\ell} > 1$ , such as  $\mathcal{O}_2$  and  $\mathcal{O}_4$  in the example. The system associated with  $\mathcal{O}_{\ell}$  is

$$F_{\alpha^r s_\ell} u_{\alpha^r s_\ell} = \lambda u_{\alpha^{r+1} s_\ell}, \quad 0 \le r \le r_\ell - 1, \quad \text{where} \quad \alpha^{r_\ell} = 1,$$

which is analogous to (ii), where  $s_{\ell} = 1$ ,  $\alpha = 3$  and k = 10. Since  $\lambda \neq 0$ , this is equivalent to

$$u_{\alpha^{r+1}s_{\ell}} = \frac{1}{\lambda} F_{\alpha^{r}s_{\ell}} u_{\alpha^{r}s_{\ell}}, \quad 0 \le r_{\ell} - 1,$$

which is analogous to (31). Therefore

$$u_{\alpha^{r+1}s_{\ell}} = \frac{1}{\lambda^{r+1}} G_{\alpha^{r+1}s_{\ell}} u_{s_{\ell}}, \quad 0 \le r \le r_{\ell} - 1,$$
(37)

where

$$G_{\alpha^{r+1}s_{\ell}} = F_{\alpha^r s_{\ell}} \cdots F_{s_{\ell}}, \quad 0 \le r \le r_{\ell} - 1,$$

which is analogous to (32) and (33). In particular, setting  $r = r_{\ell} - 1$  and noting that  $\alpha^{r_{\ell}} s_{\ell} = s_{\ell}$  yields

$$u_{s_{\ell}} = \frac{1}{\lambda^{r_{\ell}}} G_{s_{\ell}} u_{s_{\ell}} \quad \text{where} \quad G_{s_{\ell}} = F_{\alpha^{r_{\ell}-1} s_{\ell}} \cdots F_{s_{\ell}}.$$

Therefore, if  $(\gamma_{\ell}, u_{s_{\ell}})$  is an eigenvalue of  $G_{s_{\ell}}$ , then  $\gamma_{\ell}^{1/r_{\ell}}$  is an eigenvalue of A with associated eigenvector

$$z_{\ell} = \left(P_{s_{\ell}} + \sum_{m=1}^{r_{\ell}-1} \gamma_{\ell}^{-m/r_{\ell}} P_{\alpha^m s_{\ell}} G_{\alpha^m s_{\ell}}\right) u_{s_{\ell}},\tag{38}$$

which is analogous to (34). However, since  $\gamma^{1/r_{\ell}} e^{2\pi i r/r_{\ell}}$  are all  $r_{\ell}$ -th roots of  $\gamma$ , they are all eigenvalues of A. Replacing  $\gamma^{1/r_{\ell}}$  with  $\gamma^{1/r_{\ell}} e^{2\pi i r/4}$  in (38) yields associated eigenvectors

$$z_{r\ell} = \left( P_{s_{\ell}} + \sum_{m=1}^{r_{\ell}-1} \gamma_{\ell}^{-m/r_{\ell}} e^{-2\pi i r m/r_{\ell}} P_{\alpha^{m} s_{\ell}} G_{\alpha^{m} s_{\ell}} \right) u_{s_{\ell}}, \quad 0 \le r \le r_{\ell} - 1, \quad (39)$$

which is analogous to (35).

**Remark 8** Now we apply the preceding argument to a standard  $\alpha$ -circulant  $A = [a_{s-\alpha r}]_{r,s=0}^{k-1}$  with  $gcd(\alpha, k) = 1$ . From Remark 3,

$$A = \sum_{s=0}^{k-1} f_s \phi_{\alpha s} \phi_s^* \quad \text{with} \quad f_s = \sum_{r=0}^{k-1} a_r \zeta^{rs}, \quad 0 \le r \le k-1,$$

and  $\phi_0, \phi_1, \ldots, \phi_{k-1}$  as in (1). Then  $z = \sum_{s=0}^{k-1} u_s \phi_s$  is  $\lambda$ -eigenvector of A if and only if  $f_s u_s = \lambda u_{\alpha s}, 0 \le s \le k-1$ . Let  $\mathcal{O}_{\ell}$  be as in (36) and assume that  $f_{\alpha^r s_\ell} \ne 0, 0 \le r \le r_\ell - 1$ . Let

$$g_{\alpha^{r+1}s_{\ell}} = \prod_{q=0}^{r} f_{\alpha^{q}s_{\ell}}, \quad 0 \le r \le r_{\ell} - 1,$$

and

$$\gamma_{\ell} = g_{\alpha^{r_{\ell}} s_{\ell}} = f_{\alpha^{r_{\ell}-1} s_{\ell}} \cdots f_{s_{\ell}}.$$

From (37),

$$u_{\alpha^{r+1}s_{\ell}} = \frac{1}{\lambda^{r+1}} g_{\alpha^{r+1}s_{\ell}} u_{s_{\ell}}, \quad 0 \le r \le r_{\ell} - 2, \quad \text{and} \quad u_{\alpha^{r_{\ell}}s_{\ell}} = u_{s_{\ell}} = \lambda^{-r_{\ell}} \gamma_{\ell} u_{s_{\ell}}.$$

Therefore  $\gamma_{\ell}^{r_{\ell}} e^{2\pi i r/r_{\ell}}, 0 \leq r \leq r_{\ell} - 1$ , are eigenvalues of A. From (39),

$$z_{r\ell} = \left(\phi_{s_{\ell}} + \sum_{m=1}^{r_{\ell}-1} \gamma_{\ell}^{-m/r_{\ell}} e^{-2\pi i r m/r_{\ell}} g_{\alpha^{m} s_{\ell}} \phi_{\alpha^{m} s_{\ell}}\right), \quad 0 \le r \le r_{\ell} - 1,$$

are associated eigenvectors.

For example, let  $\alpha = k - 1$ , so  $A = [a_{s+r}]_{r,s=0}^{k-1}$ . If k = 2p then  $v_0 = f_0$ ,  $v_\ell = \sqrt{f_\ell f_{k-\ell}}, 1 \le \ell \le p - 1$ , and  $v_p = f_p$ . Hence,  $(f_0, \phi_0), (f_p, \phi_p),$ 

$$\left(\sqrt{f_{\ell}f_{k-\ell}}, \phi_{\ell} + \frac{1}{\sqrt{f_{\ell}f_{k-\ell}}}\phi_{(k-1)\ell}\right), \text{ and } \left(-\sqrt{f_{\ell}f_{k-\ell}}, \phi_{\ell} - \frac{1}{\sqrt{f_{\ell}f_{k-\ell}}}\phi_{(k-1)\ell}\right)$$

 $1 \leq \ell \leq p-1$ , are eigenpairs of A. If k = 2p+1 then  $v_0 = f_0$  and  $v_\ell = f_\ell f_{k-\ell}$ ,  $1 \leq \ell \leq q$ . Hence  $(f_0, \phi_0)$ ,

$$\left(\sqrt{f_{\ell}f_{k-\ell}}, \phi_{\ell} + \frac{1}{\sqrt{f_{\ell}f_{k-\ell}}}\phi_{(k-1)\ell}\right), \quad \text{and} \quad \left(-\sqrt{f_{\ell}f_{k-\ell}}, \phi_{\ell} - \frac{1}{\sqrt{f_{\ell}f_{k-\ell}}}\phi_{(k-1)\ell}\right)$$

 $1 \leq \ell \leq p$ , are eigenpairs of A.

The eigenvalues of A were given in [5] without the associated eigenvectors.

## 9 Acknowledgement

I thank the referee for suggestions that clarified and improved this paper and for correcting numerous typographical errors.

## References

- C. M. Ablow, J. L. Brenner, Roots and canonical forms for circulant matrices, Trans. Amer. Math. Soc. 107 (1963) 360–376.
- [2] A. L. Andrew, Eigenvectors of certain matrices, Linear Algebra Appl. 7 (1973) 151–162.
- [3] A. Ben-Israel and T.N.E. Greville, Generalized Inverses: Theory and Applications, John Wiley and Sons, New York, 1974.
- [4] E. C. Boman, The Moore-Penrose pseudoinverse of an arbitrary, square, k-circulant matrix. Lin. Multilin. Alg. 50 (2002) 175-179.
- [5] S. Charmonman, R. S. Julius, Explicit inverses and condition numbers of certain circulants, Math. Comput. 102 (1968) 428–430.
- [6] P. J. Davis, Cyclic transformations of polygons and the generalized inverse, Canad. J. Math. 29 (1977) 756-770.
- [7] R. M. Gray, Toeplitz and circulant matrices, Foundations and Trends in Communications and Information Theory 2 (2006) 155-239.
- [8] S. Serra Capizzano, A Korovkin-type theory for finite Toeplitz operators via matrix algebras, Numerische Mathematik 82 (1999) 117-142.
- [9] S. Serra Capizzano, A Korovkin based approximation of multilevel Toeplitz matrices (with rectangular unstructured blocks) via multilevel trigonometric matrix spaces, SIAM Journal on Numerical Analysis 36 (1999) 1831-1857.
- [10] W. T. Stallings, T. L. Boullion, The pseudoinverse of an *r*-circulant matrix, Proc. Amer. Math. Soc. 34 (1972) 385–388.